

Status of constraints on SUSY

A. Freitas

University of Pittsburgh

- 1. SUSY models**
- 2. High energy colliders**
- 3. Electroweak precision data**
- 4. Flavor and Higgs physics**
- 5. Cosmology**

Status of constraints on SUSY

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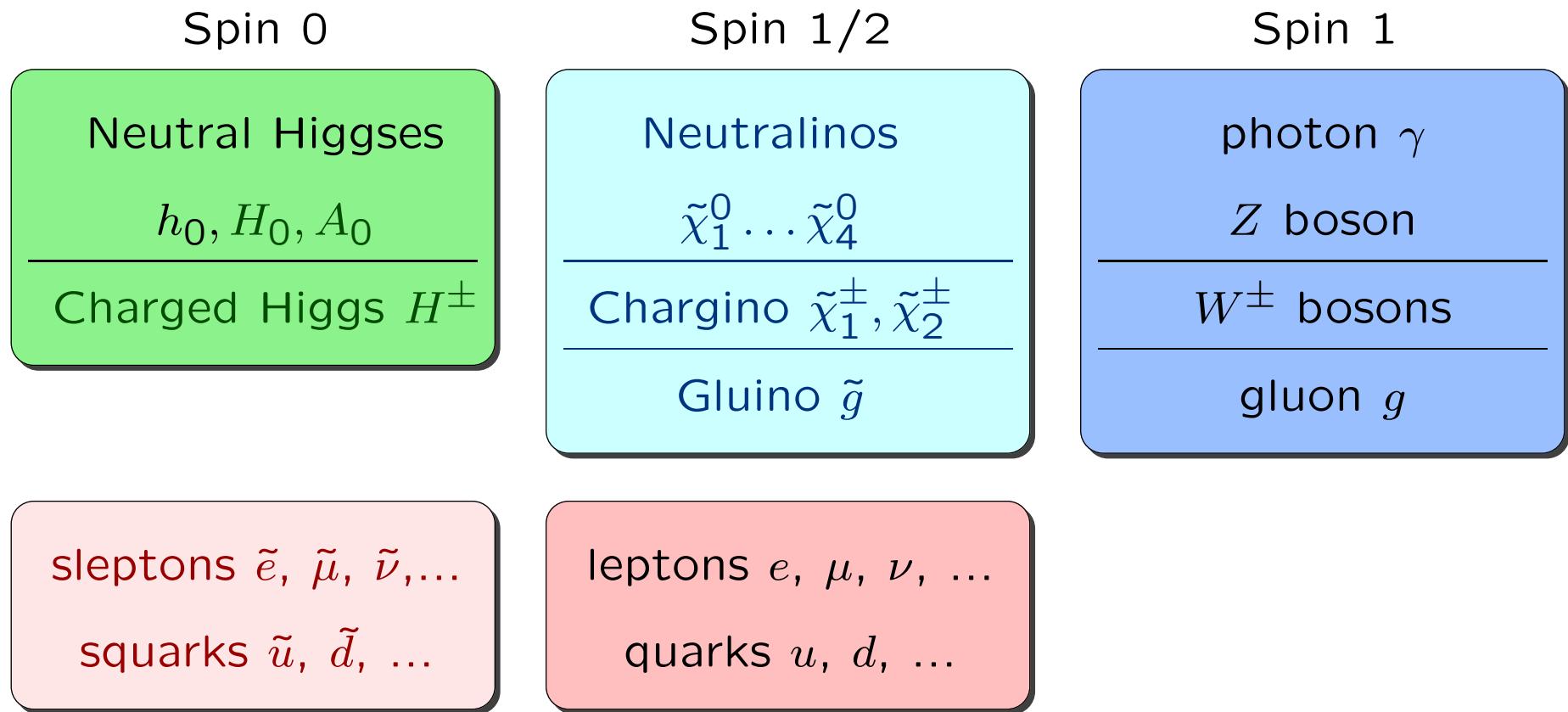
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Apologies for the very incomplete references!

Introduction: SUSY models

“Minimal MSSM”:



- R-parity conservation: $R_p = \pm 1$ for SM/SUSY particles
- CP conservation
- Minimal flavor violation: CKM matrix is only source of flavor violation

SUSY breaking

mSUGRA/CMSSM:

- Gaugino masses, scalar masses and triple-scalar couplings (A-terms) unify at $M_{\text{GUT}} \sim 10^{16}$ GeV
- Lightest SUSY particle (LSP) is $\tilde{\chi}_1^0$, $m_{\tilde{\chi}_1^0} \sim \mathcal{O}(100)$ GeV

GMSB:

- Gaugino masses unify at $M_{\text{GUT}} \sim 10^{16}$ GeV
- Scalar masses do not unify
- Triple-scalar couplings ≈ 0 at $\Lambda_{\text{mess}} \sim 100$ TeV
- LSP is gravitino $m_{\tilde{G}} \sim 100$ eV ... 1 GeV

MSSM with additional singlet

■ μ -problem in MSSM:

Superpotential has term $W_{\text{MSSM}} \supset -\mu \hat{H}_1 \cdot \hat{H}_2$

→ Why is $\mu \ll M_{\text{Planck}}$?

■ Solve μ -problem with new singlet superfield S :

$$W = \lambda \hat{S} \hat{H}_1 \cdot \hat{H}_2 + \kappa \hat{S}^3 + m_S \hat{S}^2 + \tau_S \hat{S} + \text{Yukawa terms}$$

Effective μ -term through VEV of S : $\mu_{\text{eff}} = -\lambda \langle S \rangle$

- a) NMSSM: \mathbb{Z}_3 global symmetry $\rightarrow m_S = 0, \tau_S = 0$
- b) nMSSM: $\mathbb{Z}_{5,7}$ global symmetry $\rightarrow m_S = 0, \kappa = 0, \tau_S = \mathcal{O}(\text{TeV})$
- c) UMSSM: $U(1)$ gauge symmetry $\rightarrow m_S = 0, \kappa = 0, \tau_S = 0$
(has also Z' gauge boson)

SUSY constraints from high energy colliders

Scenario-independent results from LEP2:

Sparticle	lower limit [GeV]	Sparticle	lower limit [GeV]
$\tilde{\chi}_2^0$	62.4	$\tilde{\nu}$	94.0
$\tilde{\chi}_3^0$	99.9	\tilde{e}_L	107.0
$\tilde{\chi}_4^0$	116.0	$\tilde{\mu}_R$	91.0
$\tilde{\chi}_1^\pm$	94.0	$\tilde{\tau}_1$	81.9
Some of the limits require $m_{\tilde{X}} - m_{LSP} > 3\text{--}15 \text{ GeV}$		\tilde{q}	97.0
		\tilde{t}_1	92.6

Constraints from Tevatron:

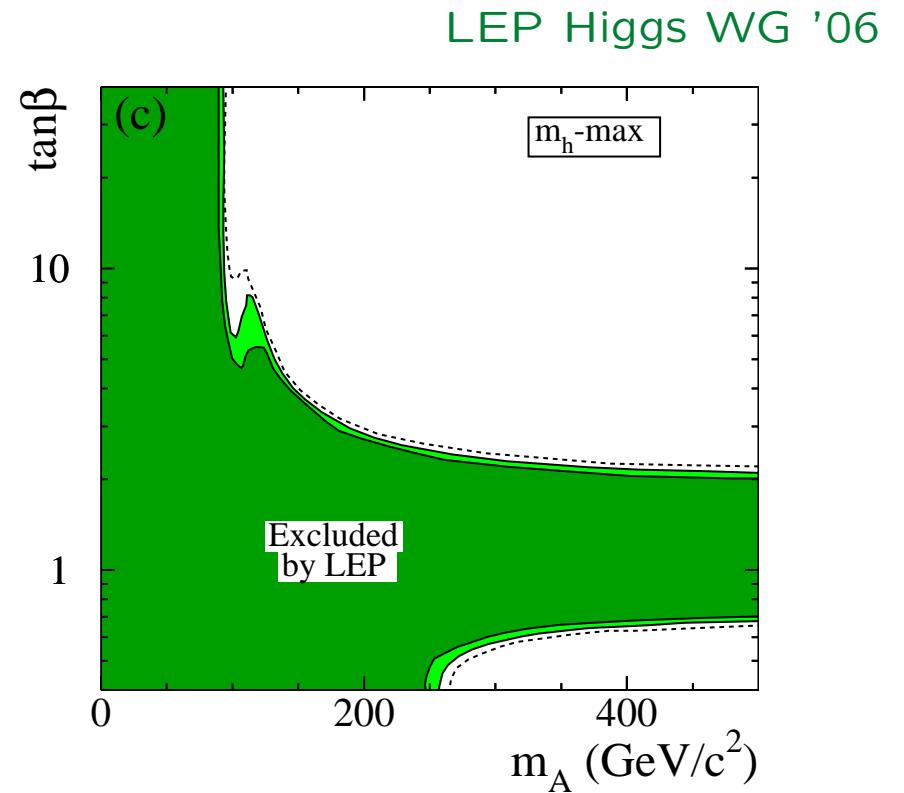
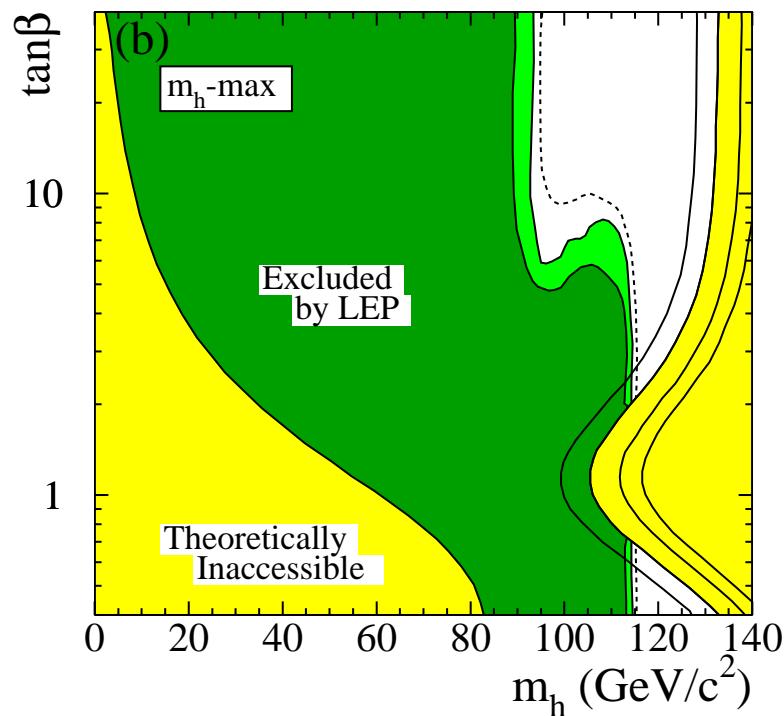
CMSSM/mSUGRA	more general
\tilde{g}	308
\tilde{q}	380
$\tilde{\chi}_1^\pm$	140

more in following talks...

Higgs searches at LEP2

LEP searched for Higgs production in $e^+e^- \rightarrow hZ$

$$\rightarrow m_h^{\text{SM}} > 114.4 \text{ GeV}$$



$$m_h \lesssim M_Z^2 \cos^2 2\beta + \frac{3m_t^4}{2\pi^2 v^2} \left[\log \frac{m_t^2}{m_{\tilde{t}}} + \frac{X_t^2}{m_{\tilde{t}}^2} + \dots \right]$$

$$X_t = A_t - \frac{\mu}{\tan\beta}$$

Higgs searches at LEP2

MSSM parameter space has to fulfill at least one of the following:

- $\tan \beta \gg 1$ because of $M_Z^2 \cos^2 2\beta$ term
- Heavy stops $m_{\tilde{t}}^2 = m_{\tilde{t}_1} m_{\tilde{t}_2} \gtrsim (1 \text{ TeV})^2$
- Large stop mixing because of $X_t^2/m_{\tilde{t}}^2$ term

Constraints on models with extra singlets (NMSSM) or gauge groups are relaxed due to new tree-level contributions

Batra, Delgado, Kaplan, Tait '04

Electroweak precision data

SUSY loop corrections affect predictions for electroweak precision observables

Current state of the art:

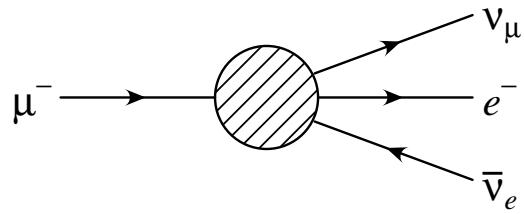
- Complete MSSM one-loop corrections Chankowski et al. '94
- Leading two-loop corrections $\mathcal{O}(\alpha\alpha_s)$ and $\mathcal{O}(\alpha y_{t,b}^2)$ Djouadi et al. '98
Haestier et al. '05
- Only partial one-loop results for NMSSM

Important quantities:

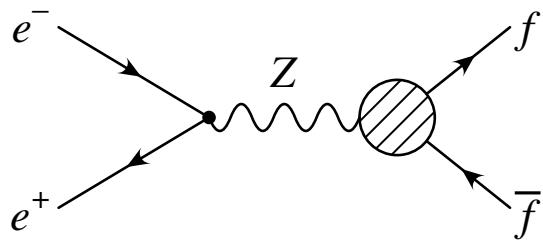
- W-boson mass M_W
- Effective weak mixing angle of Z-boson $\sin \theta_{\text{eff}} = \frac{1}{4} \left(1 - \text{Re} \frac{v_{\text{eff}}}{a_{\text{eff}}} \right)$
- Z-boson width Γ_Z and total cross section $\sigma[e^+e^- \rightarrow Z \rightarrow f\bar{f}]$
- Muon anomalous magnetic moment $a_\mu = (g_\mu - 2)/2$

Electroweak precision data

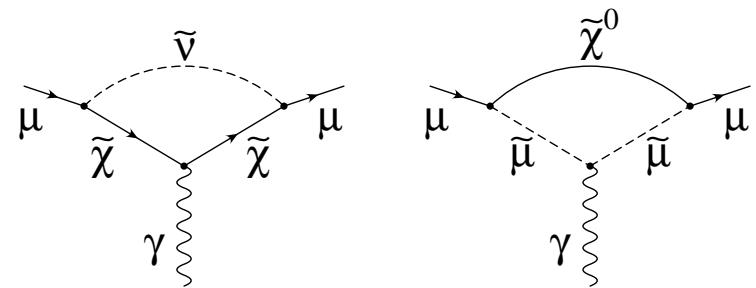
- W-boson mass M_W (from muon decay)



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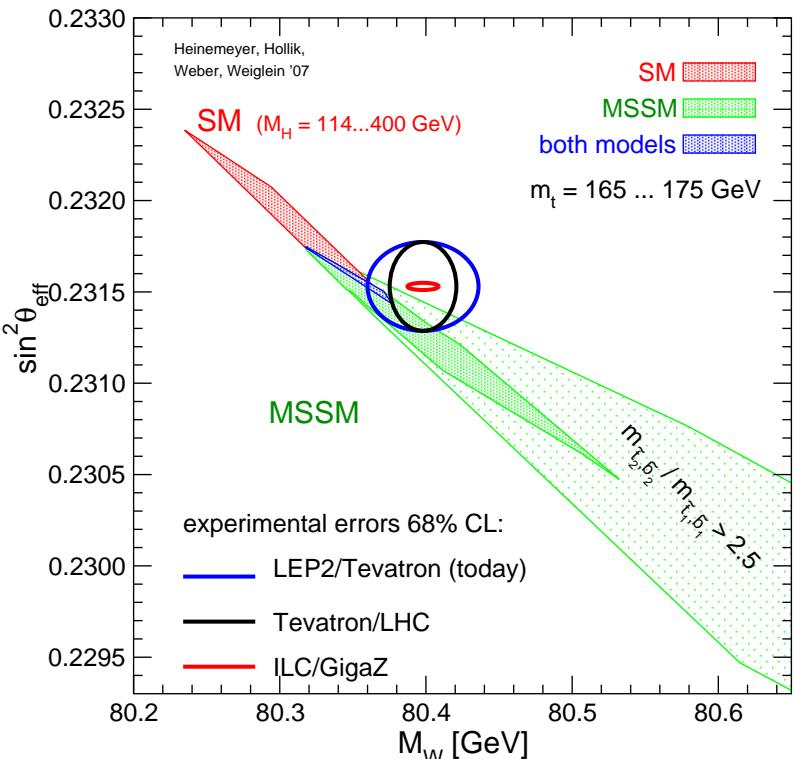
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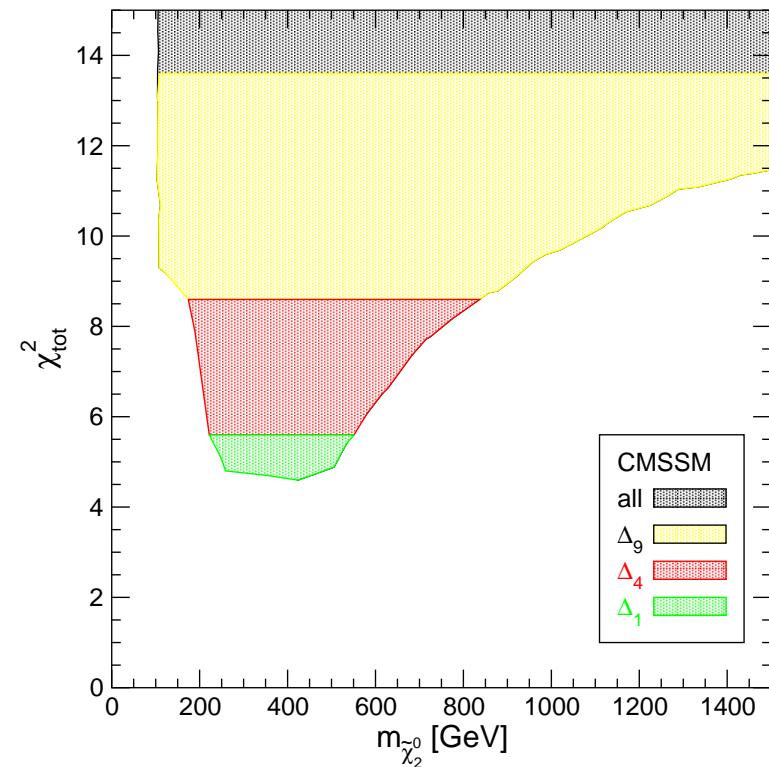
Electroweak precision data

EWPO prefer relatively light sleptons and gauginos:

- Best fit to EWPO in SM gives $m_h \approx 87$ GeV
Can be pushed to $m_h > 100$ GeV by SUSY effects
- SUSY can account for 3.3σ discrepancy
 $a_\mu^{\text{exp}} - a_\mu^{\text{theo}} = (27.5 \pm 8.4) \times 10^{-10}$



Heinemeyer, Hollik, Weber, Weiglein '07

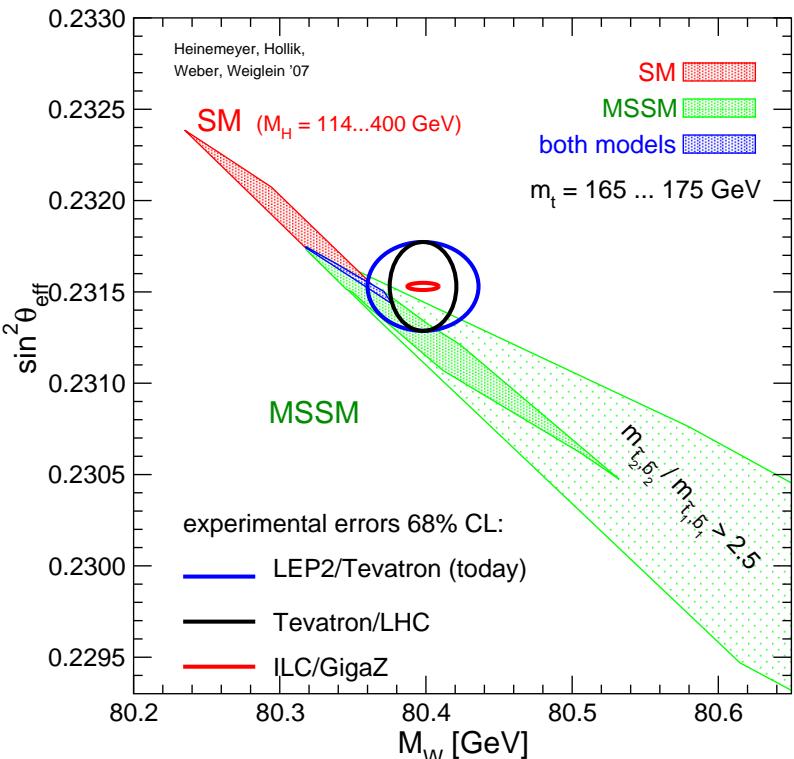


Heinemeyer, Miao, Su, Weiglein '08

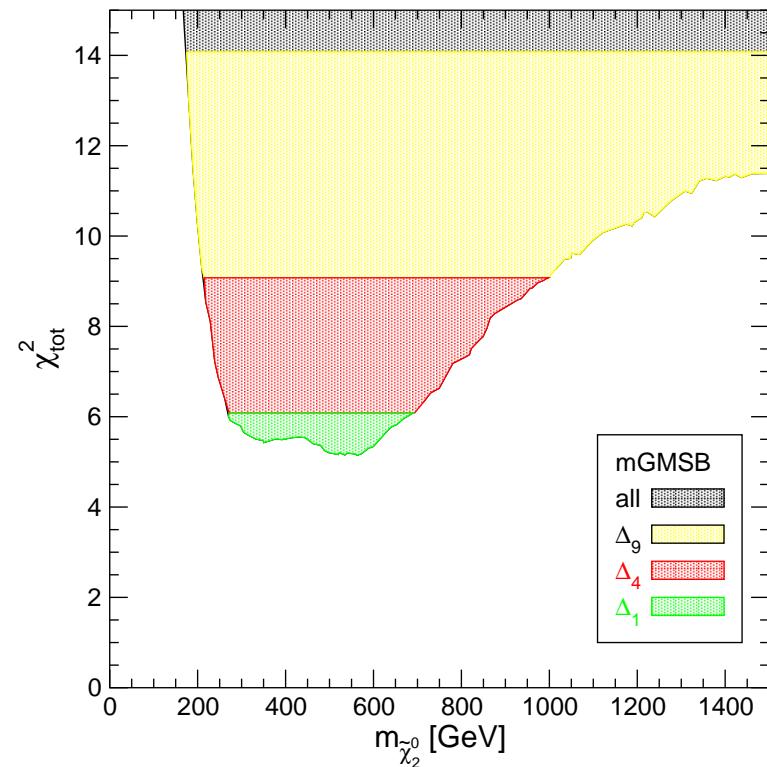
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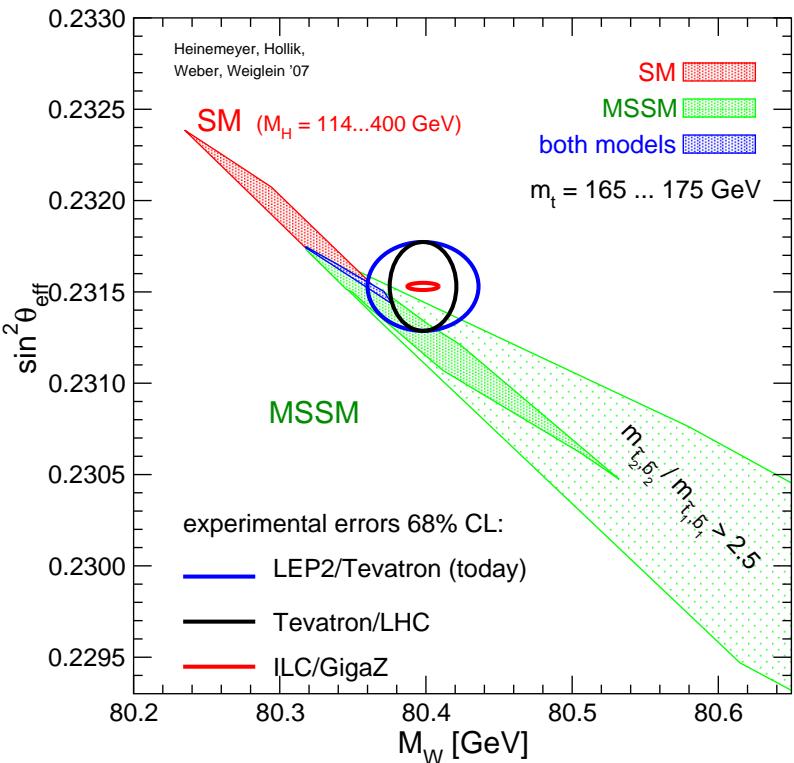


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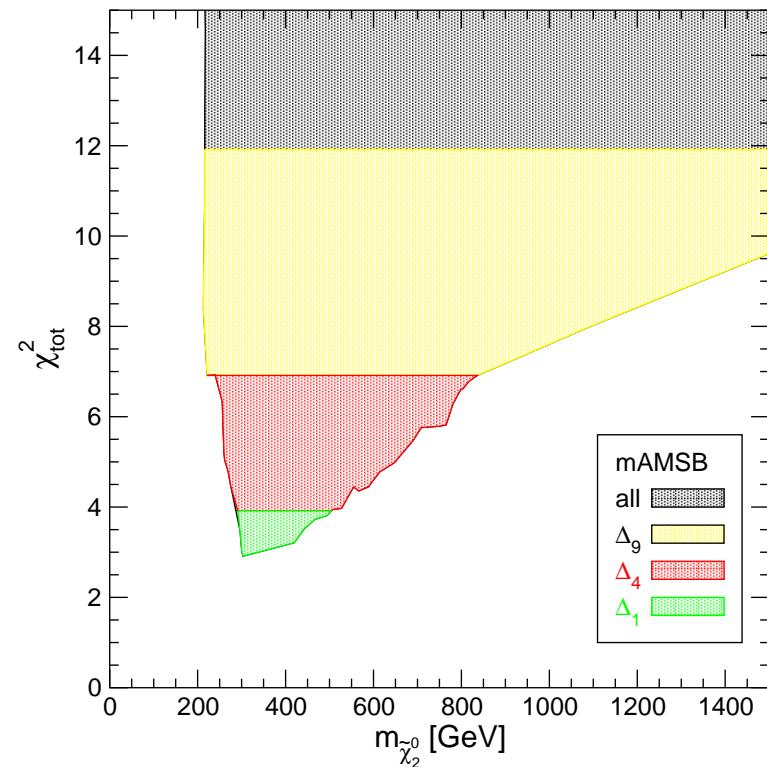
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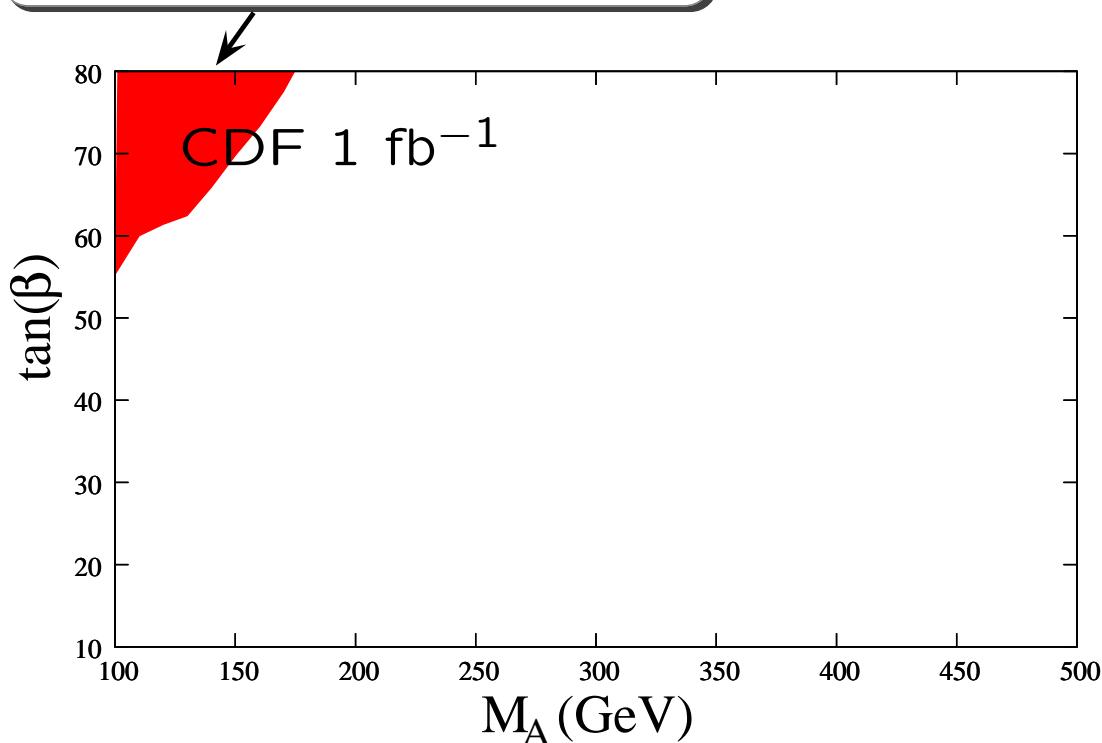
Flavor and Higgs physics

For large $\tan \beta$ enhancement for

Carena, Menon, Wagner '07

- production of non-standard Higgs bosons
- effects in B physics observables

$$\sigma[pp \rightarrow A \rightarrow \tau\tau] \sim \tan^2 \beta$$



$$\begin{aligned}\text{BR}[B_s \rightarrow \mu\mu] &\sim \tan^6 \beta \\ \text{BR}[B_u \rightarrow \tau\nu] &\sim 1 - \frac{\tan^2 \beta}{M_A^2}\end{aligned}$$

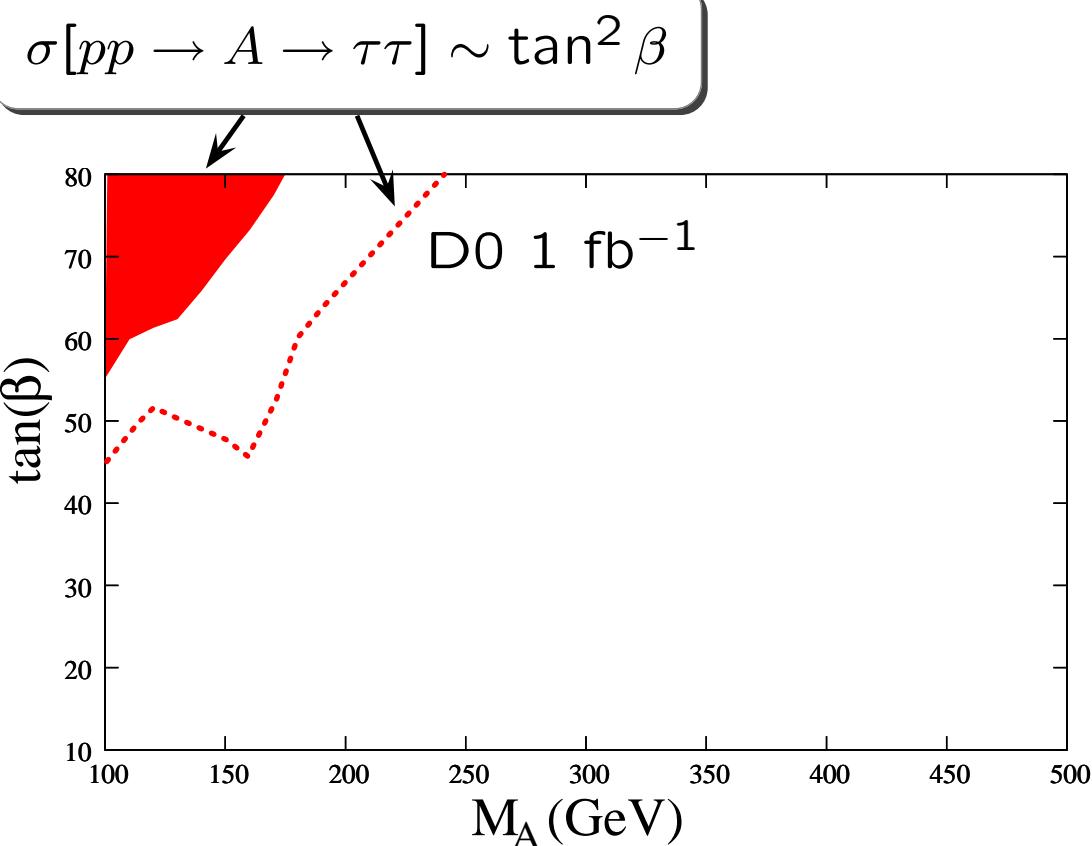
$$\begin{aligned}\text{BR}[b \rightarrow s\gamma] &\sim 1 + A \tan \beta \\ &+ B \tan \beta / M_A^2\end{aligned}$$

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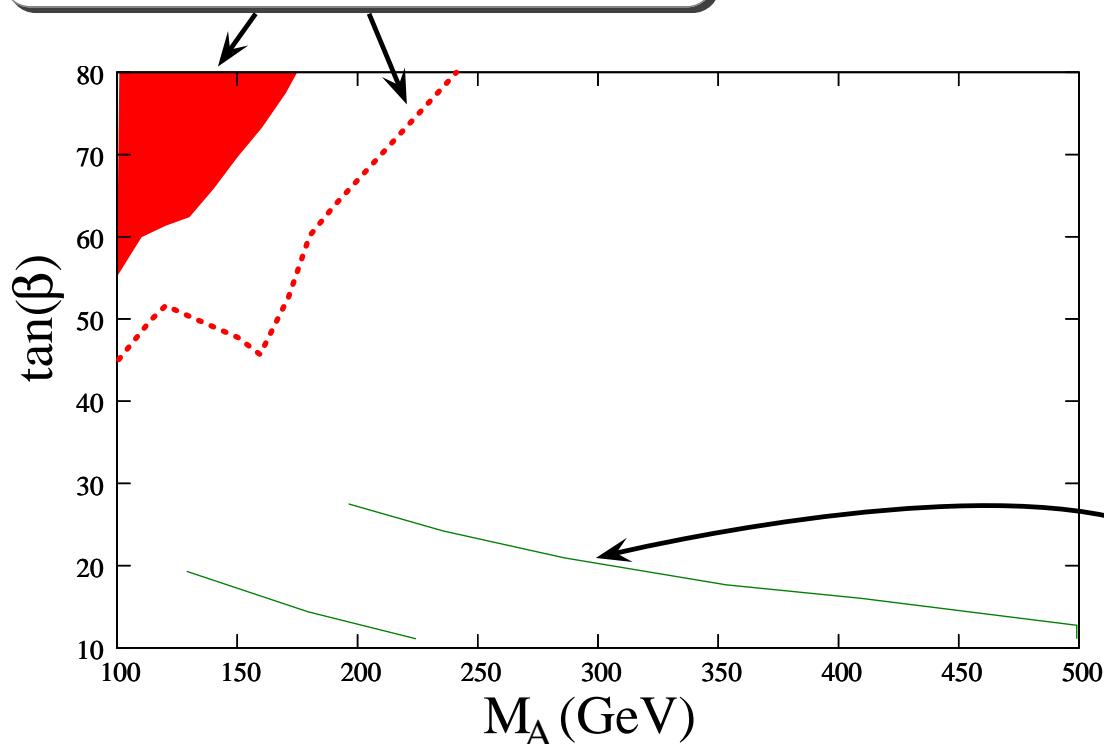
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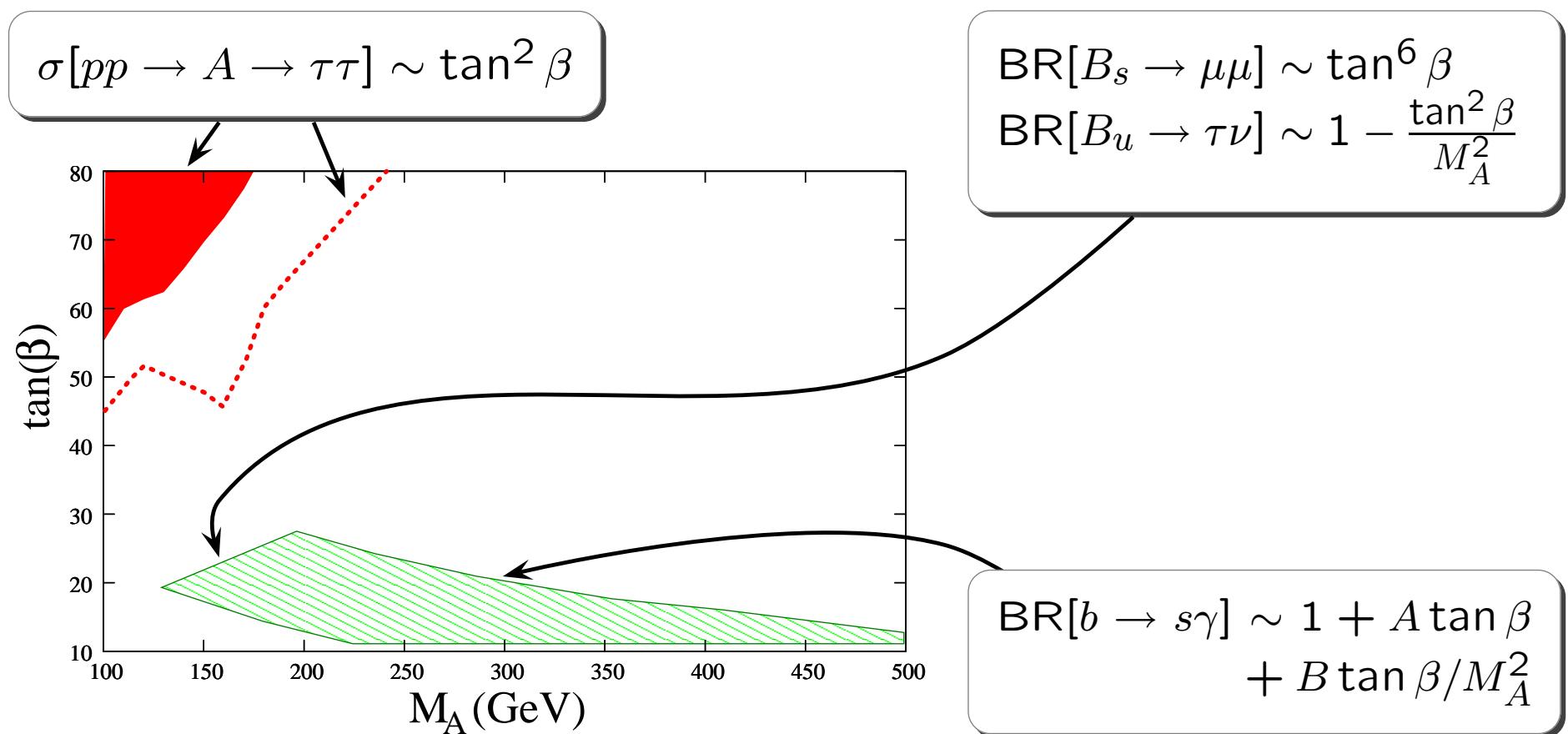
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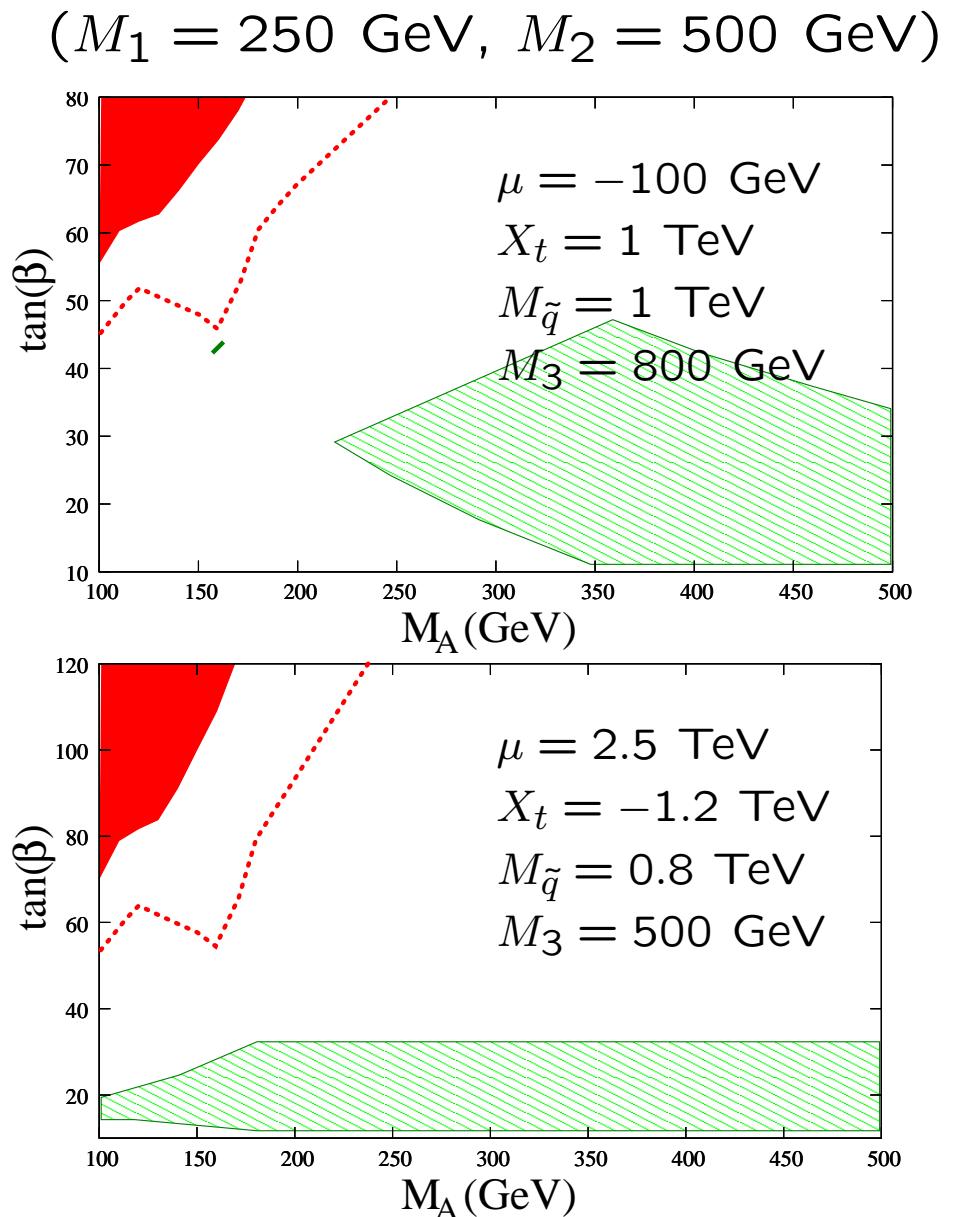
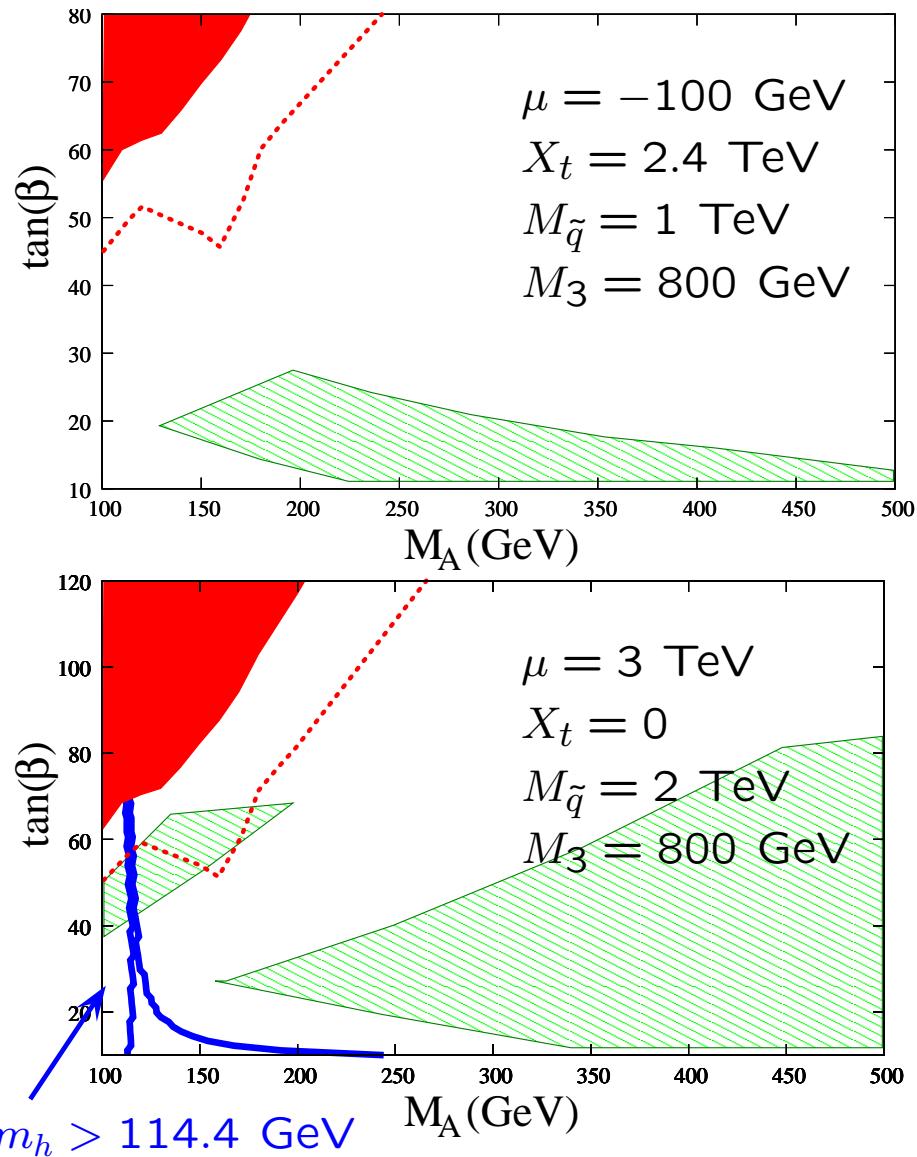
- production of non-standard Higgs bosons
- effects in B physics observables



Flavor and Higgs physics

Carena, Menon, Wagner '07

Dependence on parameters:



Flavor and Higgs physics: General picture

- Robust constraints on region of **large** $\tan\beta$ and **small** M_A from A_0 searches, $B_s \rightarrow \mu\mu$ and $B_u \rightarrow \tau\nu$
- **Large** $\tan\beta$ is also constrained for $M_A \gtrsim 250$ GeV, but depends on other parameters

Cosmology: dark matter

$$\Omega_{\text{DM}} h^2 = 0.110 \pm 0.006$$

WMAP '08

If R-parity is conserved, the LSP is stable and a **dark matter** candidate
→ Must be neutral and weakly interacting

Relic density calculation depends
on many model parameters

■ Lightest neutralino $\tilde{\chi}_1^0$:

a) $\tilde{\chi}_1^0$ is mainly bino

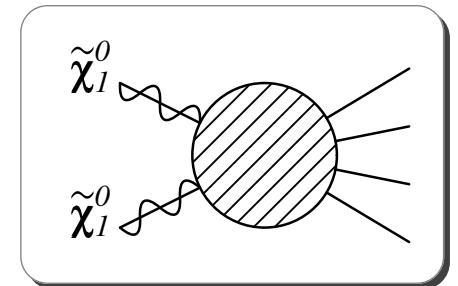
a1) $m_{\tilde{\chi}_1^0} < 100 \text{ GeV}$, $m_{\tilde{l}} \approx 100 \text{ GeV}$ (“bulk”)

a2) $m_{\tilde{X}} - m_{\tilde{\chi}_1^0} \ll m_{\tilde{\chi}_1^0}$ (co-annihilation)

a3) $2m_{\tilde{\chi}_1^0} \approx M_Z, m_h, m_A$ (resonant annihilation)

b) $\tilde{\chi}_1^0$ is mixed bino/wino/higgsino

$m_{\tilde{\chi}_1^0} \sim \mathcal{O}(\text{few } 100 \text{ GeV})$ is natural



Cosmology: dark matter

■ Sneutrino $\tilde{\nu}$:

L-Sneutrinos $\tilde{\nu}_L$ cannot be dominant source of DM
→ too large direct detection cross section

Dominantly R-sneutrinos $\tilde{\nu}_R$ can give correct Ω_{DM} in agreement with other constraints for $10 \text{ GeV} \lesssim m_{\tilde{\nu}_R} \lesssim 1 \text{ TeV}$ Arina, Fornengo '07

■ Gravitino \tilde{G} :

Steffen '06

a) \tilde{G} DM from NLSP decays

Constraints for NLSP $\tilde{\tau}$: $m_{\tilde{\tau}} \gtrsim 0.5 \text{ TeV}$ to avoid hot DM

$1 \text{ GeV} \lesssim m_{\tilde{G}} \lesssim 700 \text{ GeV}$ due to Ω_{DM}

b) Thermal \tilde{G} production

b1) Production from thermal equilibrium gives too large Ω_{DM}

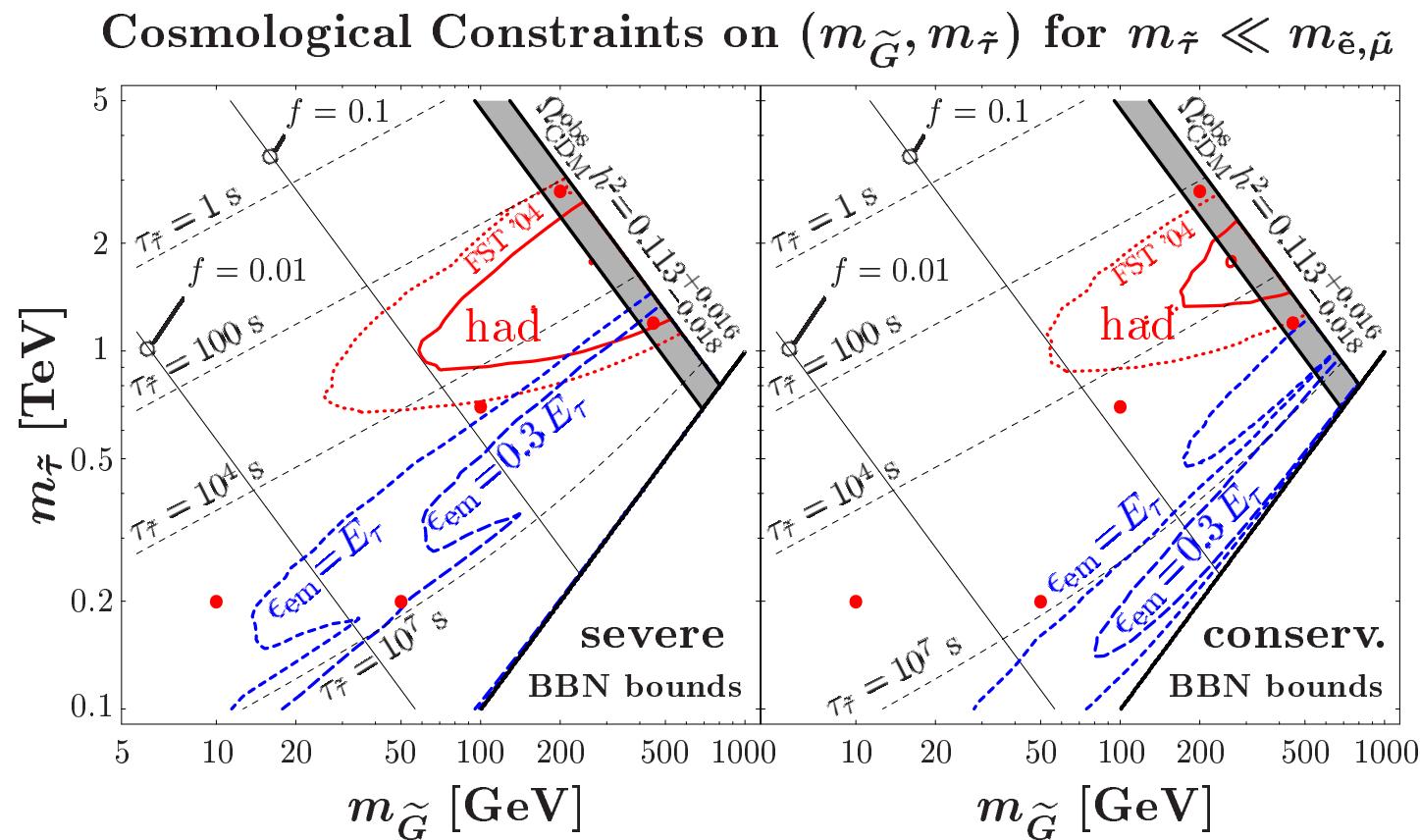
→ $T_R \ll T_{\text{eq}}$

b2) Non-equilibrium production viable for $1 \text{ keV} \lesssim m_{\tilde{G}} \lesssim 1 \text{ TeV}$, depending on T_R

Cosmology: BBN

For gravitino LSP the energy released in NLSP decays can be problematic for Big-Bang **nucleosynthesis**

For $\tau_{\text{NLSP}} \gtrsim 100$ s, hadro- and photon-dissociation can destroy light elements



Steffen '06

Cosmology: Baryogenesis

- **Leptogenesis:** Particle-Antiparticle asymmetry generated by decay of long-lived heavy (s)neutrinos
→ No constraint on low-energy spectrum

- **Electroweak baryogenesis:**

- a) Electroweak phase transition needs to be strongly **first order**

In **MSSM**: Correction to Higgs effective potential from stops

$$\rightarrow m_{\tilde{t}_1} < 140 \text{ GeV}, m_{\tilde{t}_2} > 3 \text{ TeV}$$

Carena, Quirós, Wagner '97,98
Laine, Losada, Rummukainen '98,00
Cirigliano, Profumo, Ramsey-Musolf '06

In **NMSSM**: Contributions from (tree-level) Higgs-singlet coupling
Pietroni '93; Davies, Froggatt, Moorhouse '96

- b) **CP-violating currents** active during electroweak phase transition

In **MSSM**: Complex phase in neutralino/chargino sector

$$\rightarrow \text{EDM constraints require } m_{\tilde{e}}, m_{\tilde{q}} \gtrsim 10 \text{ TeV}$$

In **NMSSM**: CP phase can also be in Higgs sector
(weaker constraints)

Huber et al. '06

Bino-like $\tilde{\chi}_1^0$ is not constrained by collider data

Constraints from astrophysics/cosmology:

- $m_{\tilde{\chi}_1^0} > 3 \text{ GeV}$ to avoid DM overproduction (R-parity conservation)
- $m_{\tilde{\chi}_1^0} > 200 \text{ MeV}$ from supernova cooling ($m_{\tilde{e}} \lesssim 500 \text{ GeV}$)
→ No constraint for $m_{\tilde{e}} \gtrsim 1200 \text{ GeV}$
- $m_{\tilde{\chi}_1^0} \gtrsim 1 \text{ keV}$ or $m_{\tilde{\chi}_1^0} \lesssim 1 \text{ eV}$ due to structure formation

Constraints on $\tilde{\chi}_1^0$ are very weak

More general SUSY breaking

Flavor violation

Sfermion mass parameters can introduce new sources of flavor violation (FV)

→ Strongly constrained by K^0, D^0 mixing, B decays,
 $\mu \rightarrow e\gamma$, $\mu \rightarrow e$ conversion, etc.

However, in 2nd and 3rd generation only, $\mathcal{O}(M_{\text{SUSY}})$ FV terms are still allowed

Ciuchini et al. '07

CP violation

Strongly constrained for gauginos and 1st generation sfermions from electron and neutron electric dipole moments (EDMs)

Sizeable CP violation allowed in Higgs sector and 3rd generation sfermions

R-parity violation

$$W_{R_p\text{MSSM}} = W_{\text{MSSM}} + \frac{1}{2} \lambda_{ijk} L_i \cdot L_j E_k^c + \frac{1}{2} \lambda'_{ijk} L_i \cdot Q_j D_k^c + \frac{1}{2} \lambda''_{ijk} U_i^c \cdot D_j^c D_k^c$$

- Product of L - and B -number violating couplings is strongly constrained by proton decay:

$$|\lambda''_{ijk} \lambda_{ijk}|, |\lambda''_{ijk} \lambda'_{ijk}| \ll 1$$

→ Could be explained by discrete symmetries

- L -number violating couplings constrained by ν masses:

$$|\lambda_{ijk}|, |\lambda'_{ijk}| \lesssim 10^{-5} \dots 0.6$$

Allanach, Dedes, Dreiner '03

- Some experimental (s)particle mass bounds get relaxed:

$$m_{\tilde{g}} > 6.3 \text{ GeV}, m_{\tilde{b}_1} > 7.5 \text{ GeV}, m_{\tilde{\tau}_1} > 11 \text{ GeV}$$

Janot '03,04

$$m_h > 82 \text{ GeV}$$

Carpenter, Kaplan, Rhee '08

Summary: rough ideas

MSSM + R_p	MSSM + \not{R}_p	NMSSM & other ext.
$m_{\tilde{l}}, m_{\tilde{q}}, m_{\tilde{\chi}_1^\pm} \gtrsim 100 \text{ GeV}, m_{H^\pm} \gtrsim 100 \text{ GeV}$		
$m_{\tilde{g}} \gtrsim 150 \text{ GeV}$	$m_{\tilde{g}} \gtrsim 6.3 \text{ GeV}$	$m_{\tilde{g}} \gtrsim 150 \text{ GeV}$
$m_{\tilde{\tau}_1} \gtrsim 82 \text{ GeV}$	$m_{\tilde{\tau}_1} \gtrsim 11 \text{ GeV}$	$m_{\tilde{\tau}_1} \gtrsim 82 \text{ GeV}$
$m_{\tilde{\mu}}, m_{\tilde{\chi}_2^0} < 1 \text{ TeV}$		
$m_{\tilde{\chi}_1^0} > 3 \text{ GeV}$	—	?
$\tan \beta \gtrsim 3$	—	—
$\sqrt{m_{\tilde{t}_1} m_{\tilde{t}_2}} \gtrsim 1 \text{ TeV}$ and/or large A_t	$(\sqrt{m_{\tilde{t}_1} m_{\tilde{t}_2}} \gtrsim 300 \text{ GeV})$	
$M_A / \tan \beta \gtrsim 3 \text{ GeV}$	—	